# A NOTE ON SPECIAL CURVES IN $E_{1}^{4}$ 

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#### Abstract

This article has consisted of a part of doctorate thesis by Süha Yılmaz [7]. Firstly, Frenet formulas are given in $E_{1}^{4}$. Later, characterizations of regular and inclined curves are studied in $E_{1}^{4}$. It has been given that a space-like curve is an inclined curve if and only if the expression $$
\frac{\kappa}{\tau}=A \cosh \left(\int_{0}^{s} \sigma d s\right)+B \sinh \left(\int_{0}^{s} \sigma d s\right)
$$ obtained using harmonic curvature functions in $E_{1}^{4}$. In addition, it has been observed that a fifth order vectorial differential equation of position vector of a space-like curve in $E_{1}^{4}$ has been satisfied by means of Frenet formulas. Similarly, it has been denoted that a fourth order vectorial differential equation of tangent vector of a space-like curve in $E_{1}^{4}$ has been also verified using Frenet formulas. Moreover, we characterized tangent and trinormal indicatrices with one theorem. Finally, it has been denoted that if and only if spherical indicatrices of space-like curve with time-like trinormal vector are regular curves.


Keywords: Inclined curve, harmonic curvature, Frenet formulas, regular curve, space-like curve, time-like curve, spherical indicatrices.

## Özet

Bu makale, Süha Yılmaz'ın doktora tezinin bir kısmını içermektedir [7]. Öncelikle, $E_{1}^{4}$ de Frenet formülleri verilmiştir. Daha sonra, regüler ve inclined eğrilerin karakterizasyonları $E_{1}{ }^{4}$ de incelenmiştir. Ayrıca, space-like bir eğrinin inclined eğri olmasının gerek yeter şartının $E_{1}^{4}$ de harmonik fonksiyonlar kullanılarak elde edilen aşağıdaki koşulu sağlaması olduğu verilmiştir:

$$
\frac{\kappa}{\tau}=A \cosh \left(\int_{0}^{s} \sigma d s\right)+B \sinh \left(\int_{0}^{s} \sigma d s\right)
$$

İlaveten, $E_{1}^{4}$ deki space-like bir eğrinin konum vektörüne ait beşinci mertebeden vektörel diferensiyel denklemi, Frenet formülleri aracılığıyla sağlanmıştır. Benzer şekilde, $E_{1}^{4}$ deki space-like bir eğrinin teğet vektörüne ait dördüncü mertebeden vektörel diferensiyel denklemi, Frenet formülleri aracılığıyla doğrulanmıştır. Ayrıca teğet ve trinormal göstergeler bir teoremle karakterize edilmiştir. Son olarak, space-like bir eğrinin küresel göstergelerinin regüler eğri olması için gerek ve yeter koşullar verilmiştir.

Anahtar Kelimeler: Inclined eğri, harmonik eğrilik, Frenet formülleri, regüler eğri, spacelike eğri, time-like eğri, küresel göstergeler.

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## 1. INTRODUCTION and PRELIMINARIES

The classical differential geometry of curves in Euclidean 4-space and higher dimensions is studied by Șemin [5], Gluck [3], and Mağden [8]. Also the studies of curves in Minkowski spaces are seen in the works of Yılmaz [7] and Ekmekçi [1].

In this study, we give some new characterizations of special curves in $E_{1}{ }^{4}$

Let denote $E_{1}^{4}$ Minkowski-4 space, i.e., the manifold Euclidean 4-space $E^{4}$ together a flat with the Lorentzian metric $\langle$,$\rangle of signature (+,+,+,-) as$

$$
\langle X, Y\rangle=x_{1} y_{1}+x_{2} y_{2}+x_{3} y_{3}-x_{4} y_{4},
$$

where $X=\left(x_{1}, x_{2}, x_{3}, x_{4}\right)$ and $Y=\left(y_{1}, y_{2}, y_{3}, y_{4}\right)$ [4]. This metric is symmetric, bi-linear and non-degenerate one.

An arbitrary vector $a$ in $E_{1}^{4}$ can have one of three Lorentzian causal characters; it can be space-like if $\langle a, a\rangle>0$ or time-like if $\langle a, a\rangle<0$ and null (light-like ) if $\langle a, a\rangle=0$. Similary, an arbitrary curve $\alpha=\alpha(s)$ in $E_{1}^{4}$ locally be space-like, time-like or null if all of its velocity vectors $\alpha^{\prime}(s)$ are respectively, space-like, time-like or null for each $s \in I \subset E$. The vectors $X, Y$ in $E_{1}^{4}$ are said to be orthogonal if $\langle X, Y\rangle=0$. Recall that the norm of an arbitrary vector $a \in E_{1}^{4}$ is given by $\|a\|=\sqrt{|\langle a, a\rangle|}$ and that the velocity of the curve $\alpha(s)$ is given by $\left\|\alpha^{\prime}(s)\right\|$. Therefore, $\alpha(s)$ is a unit speed curve if and only if $\left\langle\alpha^{\prime}(s), \alpha^{\prime}(s)\right\rangle= \pm 1$ [6].

The Lorentzian sphere of center $m=\left(m_{1}, m_{2}, m_{3}, m_{4}\right)$ and radius $r \in E^{+}$in the space $E_{1}^{4}$ defined by $S_{1}^{3}=\left\{\alpha=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \alpha_{4}\right) \in E_{1}^{4}:\langle\alpha-m, \alpha-m\rangle=r^{2}\right\}[4]$.

Let $X=X(s)$ be a space-like or time-like curve with unit speed (arc-length parameterized) in Minkowski space time $E_{1}^{4}$. The Frenet frame of $X(s)$ which is shown by $\vec{T}, \vec{N}, \vec{B}$ and $\vec{E}$ can be defined as follows:

The vector $\vec{T}$ which is tangent to the curve $X(s)$ is obtained as
$\vec{T}=\frac{d \vec{X}}{d s}$,
and the first curvature function $\kappa$ which measures the curve deviation from straight line is defined as $\kappa=\left\|\vec{T}^{\prime}\right\|$.

The vector $\vec{N}$ which is called principal normal vector is defined as
$\vec{N}=\frac{\vec{T}^{\prime}}{\kappa}$.
The third, or the binormal vector of the curve is defined as
$\vec{B}=\frac{\vec{N}+\kappa \vec{T}}{\left\|\vec{N}^{\prime}+\kappa \vec{T}\right\|}$,
and the second curvature function $\tau$ which measures the curve deviation from the plane $\{T, N\}$ is defined as $\tau=\left\|\vec{N}^{\prime}+\kappa \vec{T}\right\|$.

The fourth, or the trinormal vector is defined as
$\vec{E}=\mu(\vec{T} \wedge \vec{N} \wedge \vec{B})$,
where the exterior product of $\vec{T}, \vec{N}$, and $\vec{B}$ is defined as
$\vec{T} \wedge \vec{N} \wedge \vec{B}=-\left|\begin{array}{cccc}\overrightarrow{e_{1}} & \overrightarrow{e_{2}} & \overrightarrow{e_{3}} & -\overrightarrow{e_{4}} \\ t_{1} & t_{2} & t_{3} & t_{4} \\ n_{1} & n_{2} & n_{3} & n_{4} \\ b_{1} & b_{2} & b_{3} & b_{4}\end{array}\right|$,
here, the vectors $\overrightarrow{e_{1}}, \overrightarrow{e_{2}}, \overrightarrow{e_{3}}$, and $\overrightarrow{e_{4}}$ are coordinate directions, and the frame vectors are as $T=\left(t_{1}, t_{2}, t_{3}\right), N=\left(n_{1}, n_{2}, n_{3}\right), B=\left(b_{1}, b_{2}, b_{3}\right)$ in $E_{1}^{4}$, and also the third curvature function $\sigma$ which measures the curve deviation from the subspace $\{T, N, B\}$ is defined as
$\sigma=\mu\left(-\vec{B} \frac{d \vec{E}}{d s}\right)=\mu\left(\vec{E} \frac{d \vec{B}}{d s}\right)$.

Definition 1.1. Let $X=X(s)$ be a space-like curve with unit speed, The Frenet formulas of $X=X(s)$ for any parameter $s \in I$ are defined as
$\frac{d \vec{T}}{d s}=\kappa \vec{N}, \frac{d \vec{N}}{d s}=-\kappa \vec{T}+\tau \vec{B}, \frac{d \vec{B}}{d s}=-\tau \vec{N}+\sigma \vec{E}, \frac{d \vec{E}}{d s}=\sigma \vec{B}$,
or its matrix form is as follows

$$
\left[\begin{array}{c}
\vec{T}^{\prime}  \tag{1.7}\\
\vec{N}^{\prime} \\
\vec{B}^{\prime} \\
\vec{E}^{\prime}
\end{array}\right]=\left[\begin{array}{cccc}
0 & \kappa & 0 & 0 \\
-\kappa & 0 & \tau & 0 \\
0 & -\tau & 0 & \sigma \\
0 & 0 & \sigma & 0
\end{array}\right] \cdot\left[\begin{array}{c}
\vec{T} \\
\vec{N} \\
\vec{B} \\
\vec{E}
\end{array}\right],
$$

where $\vec{T}, \vec{N}$, and $\vec{B}$ are space-like vectors and $\vec{E}$ timelike vector of the curve $X=X(s)$, and also $\kappa, \tau$, and $\sigma$ are, respectively, the first, second, and third curvature functions of the curve $X=X(s)$.

Theorem 1.1. Let $X=X(s)$ be an arbitrary parameterized space-like curve of class $C^{5}$ in $E_{1}^{4}$, we have

$$
\begin{align*}
& \vec{T}=\frac{\dot{\vec{X}}}{\|\dot{\vec{X}}\|}, \vec{N}=\frac{\|\dot{\vec{X}}\|^{2} \ddot{\vec{X}}-(\dot{\vec{X}} \cdot \ddot{\vec{X}}) \dot{\vec{X}}}{\|\dot{\vec{X}}\|^{2} \ddot{\vec{X}}-(\dot{\vec{X}} \cdot \ddot{\vec{X}}) \dot{\vec{X}} \|}, \vec{B}=\mu \vec{E} \wedge \vec{T} \wedge \vec{N}, \text { and } \vec{E}=\dot{\mu} \frac{\vec{T} \wedge \vec{N} \wedge \vec{X}}{\|\vec{T} \wedge \vec{N} \wedge \vec{X}\|},  \tag{1.8}\\
& \kappa=\frac{\| \| \dot{\vec{X}}\left\|^{2} \ddot{\vec{X}}-(\dot{\vec{X}} \cdot \ddot{\vec{X}}) \dot{\vec{X}}\right\|}{\|\dot{\vec{X}}\|^{4}}, \tau=\frac{\|\vec{T} \wedge \vec{N} \wedge \dddot{\vec{X}}\| \cdot\|\vec{X}\|}{\| \| \dot{\vec{X}}\left\|^{2} \ddot{\vec{X}}-(\dot{\vec{X}} \cdot \ddot{\vec{X}}) \dot{\vec{X}}\right\|}, \sigma=\frac{\vec{X}^{(I V)} \cdot \vec{E}}{\|\vec{T} \wedge \vec{N} \wedge \ddot{\vec{X}}\|}, \tag{1.9}
\end{align*}
$$

where $\cdot$ shows the derivative respect to the variable $t$.
Also, it is known that the curve $X=X(s)$ is a regular curve if

$$
\begin{equation*}
\|\dot{\vec{X}}\| \neq 0 . \tag{1.10}
\end{equation*}
$$

Definition 1.2. Let us consider the space-like curve $X=X(s)$. If we transport the tangent, principal normal, and the binormal vector fields to the center $O$ of the unit hypersphere $S_{1}^{3}$, and the trinormal vector field to the center $O$ of the unit hypersphere $H_{0}{ }^{3}$, then we obtain spherical indicatrices of the curve $X=X(s)$ [7].

Theorem 1.2. [1] Let $\alpha=\alpha(s)$ be a regular curve with curvatures $\kappa \neq 0, \tau \neq 0, \sigma \neq 0$ in $E_{1}^{4}$.

The curve $\alpha$ is an inclined curve if and only if
$H_{1}^{2}-H_{2}^{2}=$ const.,
where $H_{1}$ and $H_{2}$ are the harmonic curvatures defined

$$
\begin{equation*}
H_{1}=\frac{\kappa}{\tau} \text { and } H_{2}=\frac{H_{1}^{\prime}}{\sigma} . \tag{1.12}
\end{equation*}
$$

Theorem 1.3. Let $\alpha=\alpha(s)$ be a curve in $E_{1}^{3}, \alpha=\alpha(s)$ is an inclined curve if and only if $\frac{\kappa}{\tau}=$ const. for all $s \in I$.

## 2. MAIN RESULTS

Theorem 2.1. Let $X=X(s)$ be space-like curve with time-like trinormal vector in $L^{4}$. The curve $X=X(s)$ in an inclined curve if and only if $\frac{\kappa}{\tau}=\operatorname{Ach}\left(\int_{0}^{s} \sigma d s\right)+B s h\left(\int_{0}^{s} \sigma d s\right)$,
where $\tau \neq 0, \sigma \neq 0 ; A, B$ constant.

Proof. $(\Rightarrow)$ : Let the space-like curve $X=X(s)$ be helix. In this case, from Theorem 1.2,
$H_{1}^{2}-H_{2}^{2}=$ constant.

Differentiating (2.1) with respect to variable $s$ we obtain
$H_{1} H_{1}^{\prime}-H_{2} H_{2}^{\prime}=0$,

Similarly, differentiating deriving of
$H_{2}=\frac{H_{1}}{\sigma}$,
we obtain
$H_{2}^{\prime}=\left(\frac{1}{\sigma}\right)^{\prime} H_{1}^{\prime}+\frac{1}{\sigma} H_{1}^{\prime \prime}$.
If we use (2.2), (2.3), and (2.4) we obtain
$\sigma H_{1} \cdot H_{2}-H_{2}\left[\left(\frac{1}{\sigma}\right)^{\prime} H_{1}^{\prime}+\frac{1}{\sigma} H_{1}^{\prime \prime}\right]=0$
or
$H_{2}\left[\sigma H_{1}-\left(\frac{1}{\sigma}\right)^{\prime} H_{1}^{\prime}+\frac{1}{\sigma} H_{1}^{\prime \prime}\right]=0$,
where $H_{2} \neq 0$ and $\sigma \neq 0$. Because if $\sigma=0$ and $H_{2}=0$, then we get $H_{1}^{\prime}=0$ and $H_{1}=\frac{\kappa}{\tau}=$ const. From Theorem 1.3, this shows that the curve is an inclined curve in $L^{3}$, thus the second-order homogeneous equation with variable coefficient is obtained as follows
$-\frac{1}{\sigma^{2}} \cdot H_{1}{ }^{\prime \prime}-\frac{1}{\sigma} \cdot\left(\frac{1}{\sigma}\right)^{\prime} H_{1}{ }^{\prime}+H_{1}=0$.
If the transformation $t=\int_{0}^{s} \sigma d s$ is applied to $H_{1}{ }^{\prime}, H_{2}{ }^{\prime \prime}$, we obtain
$H_{1}^{\prime}=\frac{d H_{1}}{d s}=\frac{d H_{1}}{d t} \cdot \frac{d t}{d s}=\sigma H_{1}$,
$H_{1}^{\prime \prime}=\frac{d^{2} H_{1}}{d s^{2}}=\frac{d^{2} H_{1}}{d t^{2}} \cdot\left(\frac{d t}{d s}\right)^{2}+\frac{d H_{1}}{d t} \cdot \frac{d^{2} t}{d s^{2}}=\sigma^{2} H_{1}+\sigma^{\prime} H_{1}$,
where indicates the derivative with respect to variable $t$. Substituting these equations into (2.5) we find
$-H_{1}-\frac{\sigma}{\sigma^{2}} \cdot H_{1}+\frac{\sigma}{\sigma^{2}} H_{1}+H_{1}=0$ therefore we obtain the following differential equation of constant coefficient
$-\ddot{H}_{1}+H_{1}=0$.
From solution of (2.6) we get $H_{1}=c_{1} e^{t}+c_{2} e^{-t}$ or $H_{1}=c_{1}(\cosh t+\sinh t)+c_{2}(\cosh t-\sinh t)$,
if we say $c_{1}+c_{2}=A, c_{1}-c_{2}=B$ we obtain $H_{1}=A \cosh t+B \sinh t$ or
$\frac{\kappa}{\tau}=A \cosh \left(\int_{0}^{s} \sigma d s\right)+B \sinh \left(\int_{0}^{s} \sigma d s\right)$
$(\Leftarrow)$ :Let's assume that
$\frac{\kappa}{\tau}=A \cosh \left(\int_{0}^{s} \sigma d s\right)+B \sinh \left(\int_{0}^{s} \sigma d s\right)$,
differentiating (2.7) with respect to variable $s$ we get
$\frac{1}{\sigma} \frac{d}{d s} \frac{\kappa}{\tau}=A \cosh \left(\int_{0}^{s} \sigma d s\right)+B \sinh \left(\int_{0}^{s} \sigma d s\right)$,
similarly, differentiating (2.8) with respect to variable $s$ gives
$\frac{d}{d s}\left[\frac{1}{\sigma} \frac{d}{d s} \frac{\kappa}{\tau}\right]=A \cosh \left(\int_{0}^{s} \sigma d s\right)+B \sinh \left(\int_{0}^{s} \sigma d s\right)$ or $H_{2}^{\prime}=\sigma H_{1}$ and here we obtain
$H_{2}^{\prime}-\sigma H_{1}=0$,
multiplying both side of the expression (2.9) with $H_{2}=\frac{1}{\sigma} H_{1}^{\prime}$, we obtain
$H_{1} H_{1}^{\prime}-H_{2} H_{2}^{\prime}=0$,
taking the integral of both sides of (2.10), we get
$H_{1}^{2}-H_{2}^{2}=$ const., which indicates that the curve is an inclined curve in $L^{4}$ from Theorem 1.2.

Theorem 2.2. Let $X: I \rightarrow L^{4}$ be an arc length parametrized space-like curve from class $C^{4}$ such that $\chi>0, \sigma>0$ and $\sigma=$ const. Then position vector $X=X(s)$ of the curve satisfies following a vectoral differential equation of $5^{\text {th }}$ order

$$
\begin{aligned}
& \frac{\kappa^{3}}{\tau} \cdot \vec{X}^{(v)}+\left[\left(\frac{\kappa^{3}}{\tau}\right)^{\prime}+\left(\frac{1}{\tau}\right) \cdot \frac{1}{\kappa}-\frac{2 \kappa^{\prime}}{\tau \kappa^{2}}-\frac{\sigma^{2}}{\kappa^{\tau}}\right] \vec{X}^{(v)} \\
& +2\left[\left(\frac{1}{\tau}\right)^{\prime \prime} \cdot \frac{1}{\kappa}+\left(\frac{1}{\tau}\right)^{\prime} \cdot\left(\frac{1}{\kappa}\right)^{\prime}-2\left(\frac{\kappa^{\prime}}{\tau \kappa^{2}}\right)^{\prime}-\left(\frac{\sigma^{2}}{\kappa \tau}\right)^{\prime}-\frac{\kappa^{\prime \prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime}-\frac{\kappa^{\prime \prime} \kappa^{2}}{\tau}+2 \cdot \frac{\kappa^{\prime 2}}{\kappa^{3} \tau}+\frac{\kappa}{\tau}+\frac{\tau}{\kappa}-\frac{\sigma^{2} \kappa^{\prime}}{\kappa^{2} \tau}\right] \vec{X}^{\prime \prime \prime} \\
& +\left[-\frac{\kappa^{\prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime}-\frac{\kappa^{\prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime \prime}+\left(\frac{\kappa^{\prime \prime} \kappa^{2}}{\tau}\right)+2\left(\frac{\kappa^{\prime 2}}{\kappa^{3} \tau}\right)+\left(\frac{\kappa}{\tau}\right)^{\prime}+\left(\frac{\tau}{\kappa}\right)^{\prime}-\left(\frac{\sigma^{2} \kappa^{\prime}}{\tau \kappa^{2}}\right)+\kappa\left(\frac{1}{\tau}\right)^{\prime}+\frac{\kappa^{\prime}}{\tau}-\frac{\sigma^{2}}{\tau} \kappa\right] \vec{X}^{\prime \prime} \\
& +\left[\kappa^{\prime}\left(\frac{1}{\tau}\right)^{\prime}+\kappa\left(\frac{1}{\tau}\right)^{\prime \prime}+\left(\frac{\kappa^{\prime}}{\tau}\right)^{\prime}-\left(\frac{\sigma^{2} \kappa}{\tau}\right)^{\prime}\right] \vec{X}^{\prime}=0,
\end{aligned}
$$

where $\vec{T}, \vec{N}, \vec{B}$ space-like vectors, $\vec{E}$ time-like vector, $\sigma$ constant.

Proof. Suppose $\vec{T}, \vec{N}, \vec{B}$ are space-like vectors, $\vec{E}$ is time-like vector for space-like curve $X=X(s)$ in $L^{4}$ from (1.6) and (1.6) $)_{2}$ we obtain
$\vec{N}=\frac{\vec{T}}{\kappa}$
$\vec{B}=\frac{1}{\tau}\left(\kappa \vec{T}+\vec{N}^{\prime}\right)$

Substituting (2.11) in (1.6) ${ }_{3}$ we get

$$
\begin{equation*}
\vec{B}^{\prime}=-\frac{\tau}{\kappa} \vec{T}^{\prime}+\sigma \vec{E} \tag{2.13}
\end{equation*}
$$

differentiating (2.11) and substituting it into (2.12) we find
$\vec{B}=\frac{1}{\tau}\left(\kappa \vec{T}+\left(\frac{\vec{T}}{\kappa}\right)^{\prime}\right)$
taking the integral of both sides of $(1.6)_{4}$ we get

$$
\begin{equation*}
\vec{E}=\int \sigma \vec{B} d s \tag{2.15}
\end{equation*}
$$

and substituting (2.14) into (2.15) we obtain
$\vec{E}=\int \frac{\sigma}{\tau}\left[\kappa \vec{T}+\left(\frac{\vec{T}}{\kappa}\right)^{\prime}\right] d s$
and substituting (2.16) into (2.13) we find
$\vec{B}^{\prime}=-\frac{\tau}{\kappa} \vec{T}^{\prime}+\sigma^{2} \int \frac{1}{\tau}\left[\kappa \vec{T}+\left(\frac{\vec{T}^{\prime}}{\kappa}\right)\right] d s$
similarly differentiating (2.14) and using (2.17) we obtain
$\left.\left(\frac{1}{\tau}\right)\right)^{\prime}\left[\frac{\vec{T}^{\prime \prime} \kappa-\vec{T}^{\prime} \kappa^{\prime}}{\kappa^{2}}+\kappa \vec{T}\right]+\frac{1}{\tau}\left[\left(\vec{T}^{\prime \prime} \kappa-\vec{T}^{\prime} \kappa^{\prime \prime}\right) \kappa^{2}-\frac{2 \kappa \kappa^{\prime}\left(\vec{T}^{\prime \prime} \kappa-\vec{T}^{\prime} \kappa^{\prime}\right)}{\kappa^{4}}+\kappa^{\prime} \vec{T}+\kappa \vec{T}{ }^{\prime}\right]$
$+\frac{\tau}{\kappa} \vec{T}^{\prime}-\sigma^{2} \int \frac{1}{\tau}\left[\kappa \vec{T}+\left(\frac{\vec{T}}{\kappa}\right)\right] d s=0$.
Differentiating (2.18) and substituting
$\dot{\vec{X}}=\vec{T}, \ddot{\vec{X}}=\vec{T}^{\prime}, \ddot{\vec{X}}=\vec{T}^{\prime}, \vec{X}^{(I V)}=\vec{T}^{\prime \prime \prime}, \vec{X}^{(V)}=\vec{T}^{(I V)}$ into this expression we obtain

$$
\begin{aligned}
& \frac{\kappa^{3}}{\tau} \vec{X}^{(v)}+\left[\left(\frac{\kappa^{3}}{\tau}\right)^{\prime}+\left(\frac{1}{\tau}\right) \cdot \frac{1}{\kappa}-\frac{2 \kappa^{\prime}}{\tau \kappa^{2}}-\frac{\sigma^{2}}{\kappa \tau}\right] \vec{X}^{(v)} \\
& +\left[\left(\frac{1}{\tau}\right)^{\prime \prime} \cdot \frac{1}{\kappa}+\left(\frac{1}{\tau}\right)^{\prime}\left(\frac{1}{\kappa}\right)^{\prime}-2\left(\frac{\kappa^{\prime}}{\tau \kappa^{2}}\right)^{\prime}-\left(\frac{\sigma^{2}}{\kappa \tau}\right)^{\prime}-\frac{\kappa^{\prime \prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime}-\frac{\kappa^{\prime \prime} \kappa^{2}}{\tau}+2 \cdot \frac{\kappa^{\prime 2}}{\kappa^{3} \tau}+\frac{\kappa}{\tau}+\frac{\tau}{\kappa}-\frac{\sigma^{2} \kappa^{\prime}}{\kappa^{2} \tau}\right] \vec{X}^{\prime \prime \prime} \\
& +\left[-\frac{\kappa^{\prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime}-\frac{\kappa^{\prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime \prime}+\left(\frac{\kappa^{\prime \prime} \kappa^{2}}{\tau}\right)^{\prime}+2\left(\frac{\kappa^{\prime 2}}{\kappa^{3} \tau}\right)+\left(\frac{\kappa}{\tau}\right)^{\prime}+\left(\frac{\tau}{\kappa}\right)^{\prime}-\left(\frac{\sigma^{2} \kappa^{\prime}}{\tau \kappa^{2}}\right)+\kappa\left(\frac{1}{\tau}\right)^{\prime}+\frac{\kappa^{\prime}}{\tau}-\frac{\sigma^{2}}{\tau} \kappa\right] \vec{X}^{\prime \prime} \\
& +\left[\kappa^{\prime}\left(\frac{1}{\tau}\right)^{\prime}+\kappa\left(\frac{1}{\tau}\right)^{\prime \prime}+\left(\frac{\kappa^{\prime}}{\tau}\right)^{\prime}-\left(\frac{\sigma^{2} \kappa}{\tau}\right)^{\prime}\right] \vec{X}^{\prime}=0 .
\end{aligned}
$$

Denoting these coefficients of the vectoral differential equation by

$$
P(s), Q(s), R(s), K(s), L(s)
$$

Respectively, then we can rewrite this last equation as
$P(s) \vec{X}^{(v)}+Q(s) \vec{X}^{(v)}+R(s) \vec{X}^{\prime \prime \prime}+K(s) \vec{X}^{\prime \prime}+L(s) \vec{X}^{\prime}=0$,
where
$P(s)=\frac{\kappa^{3}}{\tau}$,
$Q(s)=\left(\frac{\kappa^{3}}{\tau}\right)^{\prime}+\left(\frac{1}{\tau}\right)^{\prime} \cdot \frac{1}{\kappa}-\frac{2 \kappa^{\prime}}{\tau \kappa^{2}}-\frac{\sigma^{2}}{\kappa \tau}$,
$R(s)=\left[\left(\frac{1}{\tau}\right)^{\prime \prime} \cdot \frac{1}{\kappa}+\left(\frac{1}{\tau}\right)^{\prime}\left(\frac{1}{\kappa}\right)^{\prime}-2\left(\frac{\kappa^{\prime}}{\tau \kappa^{2}}\right)^{\prime}-\left(\frac{\sigma^{2}}{\kappa \tau}\right)^{\prime}-\frac{\kappa^{\prime \prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime}-\frac{\kappa^{\prime \prime} \kappa^{2}}{\tau}+2 \frac{\kappa^{\prime 2}}{\kappa^{3} \tau}+\frac{\kappa}{\tau}+\frac{\tau}{\kappa}-\frac{\sigma^{2} \kappa^{\prime}}{\kappa^{2} \tau}\right]$,
$K(s)=\left[-\frac{\kappa^{\prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime}-\frac{\kappa^{\prime}}{\kappa^{2}}\left(\frac{1}{\tau}\right)^{\prime \prime}+\left(\frac{\kappa^{\prime \prime} \kappa^{2}}{\tau}\right)^{\prime}+2\left(\frac{\kappa^{\prime 2}}{\kappa^{3} \tau}\right)+\left(\frac{\kappa}{\tau}\right)^{\prime}+\left(\frac{\tau}{\kappa}\right)^{\prime}-\left(\frac{\sigma^{2} \kappa^{\prime}}{\tau \kappa^{2}}\right)^{\prime}+\kappa\left(\frac{1}{\tau}\right)^{\prime}+\frac{\kappa^{\prime}}{\tau}-\frac{\sigma^{2}}{\tau} \kappa\right]$,
$L(s)=\kappa^{\prime}\left(\frac{1}{\tau}\right)^{\prime}+\kappa\left(\frac{1}{\tau}\right)^{\prime \prime}+\left(\frac{\kappa^{\prime}}{\tau}\right)^{\prime}-\left(\frac{\sigma^{2} \kappa}{\tau}\right)^{\prime}$,
$\kappa$ is the first curvature, $\tau$ the second curvature, $\sigma$ the third curvature and constant.
The equation (2.19) has a solution by Chebyshev-Matrix Method [2]. If , $P(s), Q(s), R(s), K(s)$ and $L(s)$ are one-variable functions, then these functions satisfy the following conditions:
i) They must be satisfied by 5 th order differential equation .
ii) They must be formed into Taylor series.

If $\kappa \neq 0, \tau \neq 0$ in ordinary differential equation (2.19), they can be formed Maclaurin series.

In this situation, the equation (2.19) is transformed to a matrix form with ChebyshevMatrix Method. Thus, solution of the equation (2.19) can be obtained with an analytic or approximate method.

Theorem 2.3. Let $X=X(s)$ be arc length parametrized space-like curve with time-like trinormal vector from class $C^{4}$ in $L^{4}$ such that $\kappa>0, \tau>0, \sigma>0$. Then tangent vector $\vec{T}$ of $X=X(s)$ satisfies following a vectoral differential equation of $4^{\text {th }}$ order

$$
\begin{aligned}
& \vec{T}^{(v)}+\left[\left(\frac{\kappa}{\tau}\right)^{\prime \prime} \frac{\tau}{\kappa}+2\left[\ln \left|\frac{\kappa}{\tau}\right|\right]+\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{1}{\sigma}\right] \vec{T}^{\prime \prime \prime}+\left[\left[\ln \left|\frac{\kappa}{\tau}\right|\right] \cdot\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{1}{\sigma}-\frac{\sigma^{2}}{\kappa}-1\right] \vec{T}^{\prime \prime} \\
& +\left[\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{\tau^{2}}{\kappa \sigma}-\left(\frac{\tau}{\kappa}\right)\left(\frac{\tau}{\kappa}\right)^{\prime}-\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{1}{\sigma}-2\left[\ln \left|\frac{\kappa}{\tau}\right|\right]\right] \vec{T}^{\prime}+\left[-\left(\frac{\kappa}{\tau}\right)^{\prime \prime} \cdot\left(\frac{\tau}{\kappa}\right)-\left[\ln \left|\frac{\kappa}{\tau}\right|\right] \cdot\left(\frac{\sigma}{\kappa}\right) \cdot \frac{1}{\sigma}-\frac{\sigma^{2}}{\kappa}\right] \vec{T}=0
\end{aligned}
$$

Proof. Suppose $\vec{T}, \vec{N}, \vec{B}$ are space-like vectors and $\vec{E}$ is a time-like vector, taking the derivative of Frenet formulas (1.6) with respect to arc length parametrized $s_{t}$ of tangent indicatrix, we obtain
$\vec{T}^{\prime}=\vec{N}$,
$\vec{N}^{\prime}=-\vec{T}+\frac{\tau}{\kappa} \vec{B}$,
$\vec{B}=-\frac{\tau}{\kappa} \vec{N}+\frac{\sigma}{\kappa} \vec{E}$,
$\vec{E}^{\prime}=\frac{\sigma}{\kappa} \vec{B}$,

From $(2.20)_{2}$ and $(2.20)_{3}$ we get
$\vec{B}=\frac{\kappa}{\tau}\left(\vec{N}^{\prime}-\vec{T}\right)$,
$\vec{E}=\frac{\tau}{\sigma} \vec{N}+\frac{\kappa}{\sigma} \vec{B}$,
And substituting (2.21) into (2.20) ${ }_{4}$ we get
$\vec{E}^{\prime}=\frac{\sigma}{\tau}\left(\vec{T}-\vec{N}^{\prime}\right)$

Differentiating (2.21) and using $(2.20)_{3}$ we have
$\left(\frac{\kappa}{\tau}\right)^{\prime}\left(\vec{T}-\vec{N}^{\prime}\right)+\left(\frac{\kappa}{\tau}\right)\left(\overrightarrow{T^{\prime}}-\vec{N} \vec{N}^{\prime \prime}\right)=\frac{\tau}{\kappa} \vec{N}-\frac{\sigma}{\kappa} \vec{E}$
taking the first derivative of (2.21) and substituting $(2.20)_{1},(2.21),(2.22)$ and (2.23) in this
expression we obtain,

$$
\begin{align*}
& \vec{T}^{(v)}+\left[\left(\frac{\kappa}{\sigma}\right)^{\prime \prime} \frac{\tau}{\chi}+2\left[\ln \left|\frac{\kappa}{\tau}\right|\right]^{\prime}+\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{1}{\sigma}\right] \vec{T}^{\prime \prime \prime}+\left[\left[\ln \left|\frac{\kappa}{\tau}\right|\right] \cdot\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{1}{\sigma}-\frac{\sigma^{2}}{\kappa}-1\right] \vec{T}^{\prime \prime} \\
& +\left[\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{\tau^{2}}{\kappa \sigma}-\left(\frac{\tau}{\kappa}\right) \cdot\left(\frac{\tau}{\kappa}\right)^{\prime}-\left(\frac{\sigma}{\kappa}\right) \frac{1}{\sigma}-2\left[\ln \left|\frac{\kappa}{\tau}\right|\right]\right]^{\prime}  \tag{2.25}\\
& +\left[-\left(\frac{\kappa}{\tau}\right)^{\prime} \cdot\left(\frac{\tau}{\kappa}\right)-\left[\ln \left|\frac{\kappa}{\tau}\right|\right] \cdot\left(\frac{\sigma}{\kappa}\right)^{\prime} \cdot \frac{1}{\sigma}-\frac{\sigma^{2}}{\kappa}\right] \vec{T}=0
\end{align*}
$$

Theorem 2.4. Suppose $X=X(s), Y=Y(s)$ are space-like curves with time-like trinormal vector in $L^{4}$ and let the first curvature of $X=X(s)$ be constant. If trinormal indicatrix of $X=X(s)$ is tangent indicatrix of $Y=Y(s)$, then the third curvature of $Y=Y(s)$ is constant.

Proof. Let's calculate Frenet formulas of curve $Y=Y(s)$. Let $\vec{T}_{x}, \vec{N}_{x}, \vec{B}_{x}, \vec{E}_{x}, \kappa_{x}, \tau_{x}, \sigma_{x}$, and $\vec{T}_{y}, \vec{N}_{y}, \vec{B}_{y}, \vec{E}_{y}, \kappa_{y}, \tau_{y}, \sigma_{y}$ be Frenet elements of curve $X=X(s)$ and $Y=Y(s)$, respectively. Suppose $s_{y}$ be arc length parametrized of $Y=Y(s)$. Then, we can write

$$
\begin{equation*}
\vec{Y}=\int \vec{E}_{x}(s) d s \tag{2.26}
\end{equation*}
$$

differentiating both sides of (2.26) with respect to $s$, we have
$\frac{d \vec{Y}}{d s}=\frac{d Y}{d s_{y}} \cdot \frac{d s_{y}}{d s}=\vec{E}_{x}$.

Since

$$
\begin{equation*}
\frac{d \vec{Y}}{d s_{y}}=\vec{T}_{y}, \tag{2.27}
\end{equation*}
$$

we get
$\vec{T}_{y} \cdot \frac{d s_{y}}{d s}=\vec{E}_{x}$,
and using this expression, we obtain

$$
\begin{equation*}
\vec{T}_{y}=\vec{E}_{x} \text { and } \frac{d s_{y}}{d s}=1, \tag{2.29}
\end{equation*}
$$

taking the derivatives of both side of (2.29) with respect to $s$, we get $\chi_{y} \cdot \vec{N}_{y}=\sigma_{x} \cdot \vec{B}_{x}$ and here we obtain $\vec{N}_{y}=\vec{B}_{x}$ and $\chi_{y}=\sigma_{x}$, taking the derivative of both sides of expression $\vec{N}_{y}=\vec{B}_{x}$ with respect to $s$ we obtain

$$
-\kappa_{y} \cdot \vec{T}_{y}+\tau_{y} \cdot \vec{B}_{y}=-\tau_{x} \cdot \vec{N}_{x}+\sigma_{x} \cdot \vec{E}_{x}
$$

and the expression, we have

$$
\begin{equation*}
\tau_{y}=\tau_{x} \text { and } \vec{B}_{y}=-\vec{N}_{x} . \tag{2.30}
\end{equation*}
$$

If exterior product of $\vec{T}_{y} \wedge \vec{N}_{y} \wedge \vec{B}_{y}$ is formed, we obtain

$$
\begin{equation*}
\vec{E}_{y}=-\vec{T}_{x}, \tag{2.31}
\end{equation*}
$$

and taking the derivative of both side of (2.30) with respect to $s$, we get

$$
\begin{equation*}
\sigma_{y} \vec{B}_{y}=-\kappa_{x} \vec{N}_{x} . \tag{2.32}
\end{equation*}
$$

Since $\vec{B}_{y}=-\vec{N}_{x}$ from $(2.30)_{2}$ we get

$$
\begin{equation*}
\sigma_{y}=\kappa_{x} . \tag{2.33}
\end{equation*}
$$

ince $\kappa_{x}$ is constant, $\sigma_{y}$ is also found as constant by using (2.33).

Theorem 2.5. Let tangents, principal normals, binormals and trinormals indicatrices of space-
like curve $X$ be $X_{1}, X_{2}, X_{3}, X_{4}$, respectively. If spherical indicatrices of

$$
\begin{aligned}
X: I & \rightarrow L^{4} \\
s & \rightarrow X(s)
\end{aligned}
$$

satisfy following conditions, then they are regular curve
i) The curve $X_{1}$ is regular $\Leftrightarrow \kappa \neq 0$
ii) The curve $X_{2}$ is regular $\Leftrightarrow \sqrt{\kappa^{2}-\tau^{2}} \neq 0 \quad(\kappa \neq 0, \tau \neq 0)$
iii) The curve $X_{3}$ is regular $\Leftrightarrow \sqrt{\left|\tau^{2}-\sigma^{2}\right|} \neq 0 \quad(\tau \neq 0, \sigma \neq 0)$
iv) The curve $X_{4}$ is regular $\Leftrightarrow \sigma \neq 0$

Proof. From (1.10)
i) Since,

$$
X_{1}=\vec{T} \Rightarrow \frac{d X_{1}}{d s}=\kappa \vec{N},
$$

the curve $X_{1}$ is regular $\Leftrightarrow\left\|\frac{d X_{1}}{d s}\right\|=\kappa \neq 0$
ii) Similarly, since
$X_{2}=\vec{N} \Rightarrow \frac{d X_{2}}{d s}=-\kappa \vec{T}+\tau \vec{B}$,
the curve $X_{2}$ is regular $\Leftrightarrow\left\|\frac{d X_{2}}{d s}\right\|=\sqrt{\kappa^{2}+\tau^{2}} \quad(\kappa \neq 0, \tau \neq 0)$
iii) Also, since
$X_{3}=\vec{B} \Rightarrow \frac{d X_{3}}{d s}=-\tau \vec{N}+\sigma \vec{E}$,
the curve $X_{3}$ is regular $\Leftrightarrow \frac{d X_{3}}{d s}=\sqrt{\left|\tau^{2}-\sigma^{2}\right|}(\tau \neq 0, \sigma \neq 0)$
iv) Finally, since
$X_{4}=\vec{E} \Rightarrow \frac{d X_{4}}{d s}=\sigma \vec{B}$,

The curve $X_{4}$ is regular $\Leftrightarrow\left\|\frac{d X_{2}}{d s}\right\|=\sigma \neq 0$. with time-like trinormal vector.

## REFERENCES

[1] Ekmekçi,N., Lorentz manifoldları üzerinde eğilim çizgileri, Doktora Tezi , A.Ü Fen Bilimleri Enstitüsü, 1991.
[2] Fox, L., Parker, I. B., Chebyshev Polynomials in Numerical Analysis, Oxford University Press, 1968.
[3] Gluck, H., Higher Curvatures of Curves in Euclidean Space,
Proc.America.Math.Monthly, 73,699-704, 1966.
[4] O’Neill, B., Semi-Riemannian Geometry, Academic Press, NY, 1983.
[5] Şemin, F., Diferansiyel Geometri-I, İstanbul Üniversitesi Fen Fakültesi Yayınları, 1983.
[6] Turgut, A., 3-Boyutlu Minkowski Uzayında Space-Like ve Time-Like Regle Yüzeyler.
A.Ü. Fen Bilimleri Enstitüsü, Doktora Tezi, 1995.
[7] Yilmaz, S., Spherical Indicators of Curves and Characterisations some Special Curves In Four Dimensional Lorentzian Space, Dokuz Eylül Üniversitesi Fen Bilimleri Enstitüsü, Doktora Tezi, 2001.
[8] Mağden, A., $R^{4}$ Uzayında Bazı Özel Eğriler ve Karakterizasyonları, Atatürk Üniversitesi Fen Bilimleri Enstitüsü Doktora Tezi, 1990.


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